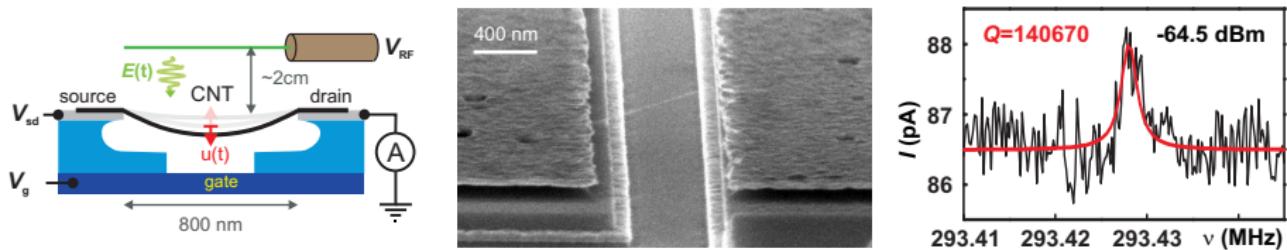


Carbon nanotubes as ultrahigh-Q electromechanical resonators at 300MHz

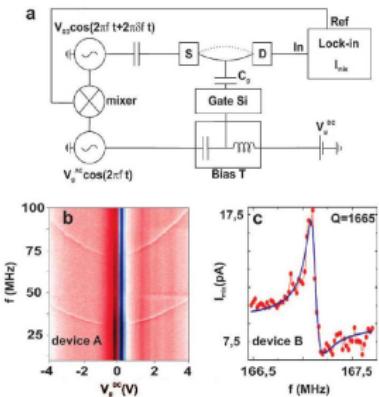
Andreas K. Hüttel*, Gary A. Steele, Benoit Witkamp, Menno Poot
Leo P. Kouwenhoven, Herre S. J. van der Zant



*Present address: Institute for Experimental and Applied Physics,
University of Regensburg, Germany

Nanotubes as beam resonators — up to now

complicated setup — even at 1K, maximally $Q \simeq 2000$

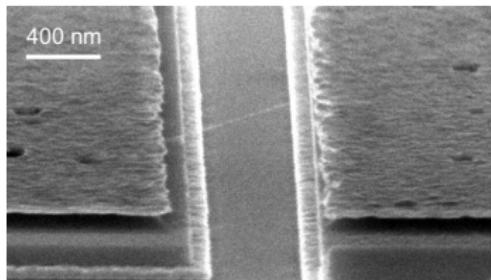
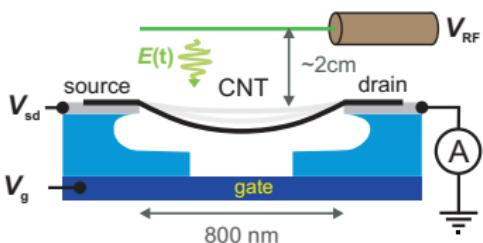


Ultrasensitive Mass Sensing with a Nanotube Electromech. Resonator
B. Lassagne, D. Garcia-Sanchez, A. Aguasca and A. Bachofeld
Nano Lett., 2008, 8 (11), pp 3735–373

- Nanotube as nonlinear circuit element
- RF downmixing at mech. resonance
- $Q \lesssim 2000$ — why?

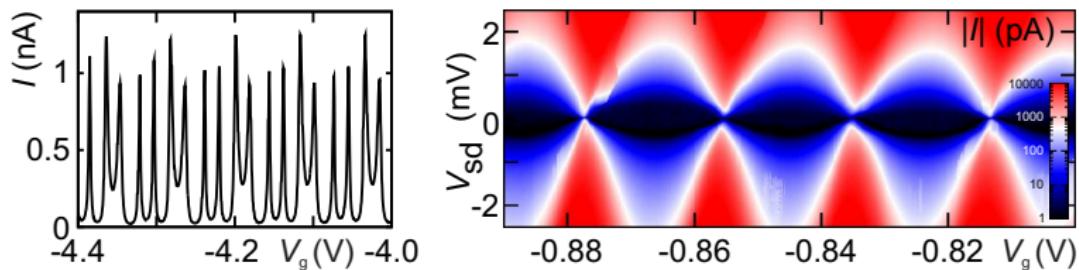
- HF cables to sample: heating, noise
- Contamination during lithography
- Clamping points?

Chip fabrication and measurement setup

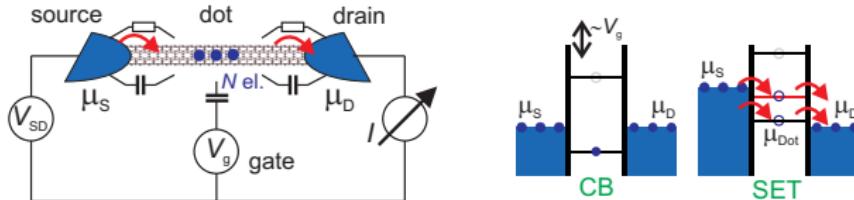


- Nanotube CVD-grown above Pt electrodes, across pre-defined trench
- Back gate connected to a gate voltage source V_g
- RF antenna suspended $\sim 2\text{cm}$ above chip
- Dilution refrigerator ($T \simeq 20\text{ mK}$)
- Only dc measurement

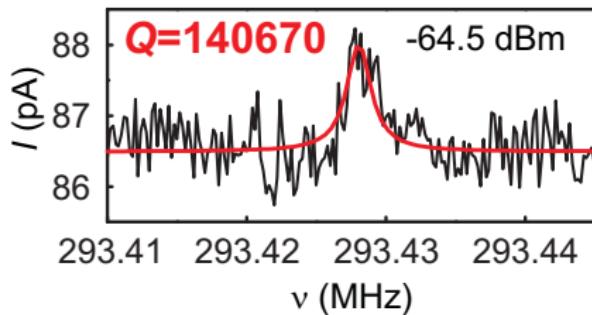
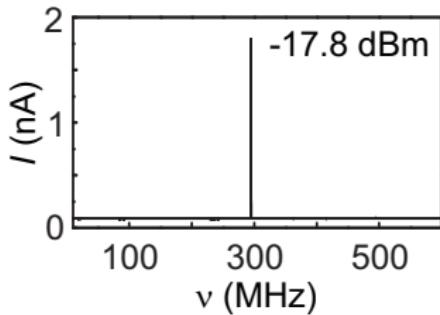
dc Coulomb blockade measurement — beautiful diamonds



highly regular quantum dot within the nanotube

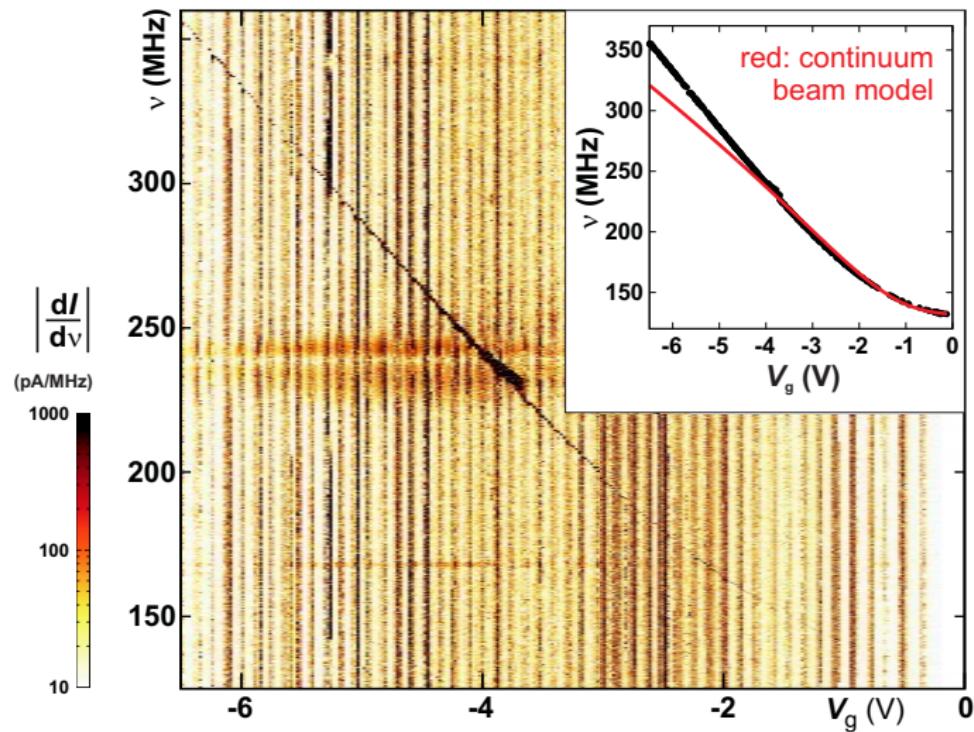


Fixed V_g and V_{SD} , sweep of RF signal frequency



- Sharp resonant structure in $I_{dc}(v)$
- Very low driving power required
- From FWHM, $Q \simeq 140000$

V_g dependence — this is really a mechanical resonance!

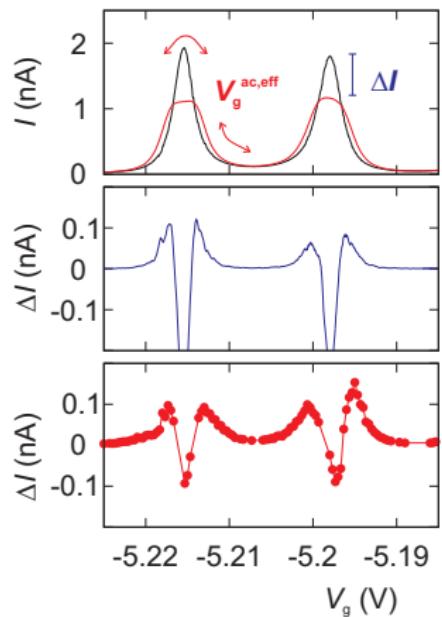


larger $|V_g| \longrightarrow$ increased tension \longrightarrow higher frequency ν

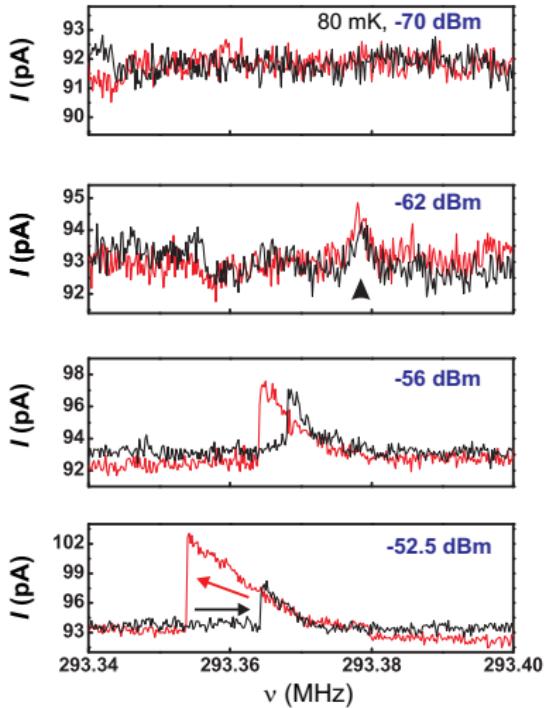
Detection mechanism — mechanically induced averaging

- at resonant driving the nanotube position oscillates
- oscillating C_g
→ fast averaging over $I(V_g)$

- black line: dc measurement $I(V_g)$
- red line: this numerically averaged
- blue: difference, effect of averaging
- red points: measured peak amplitude in $I(v)$, for different values of V_g

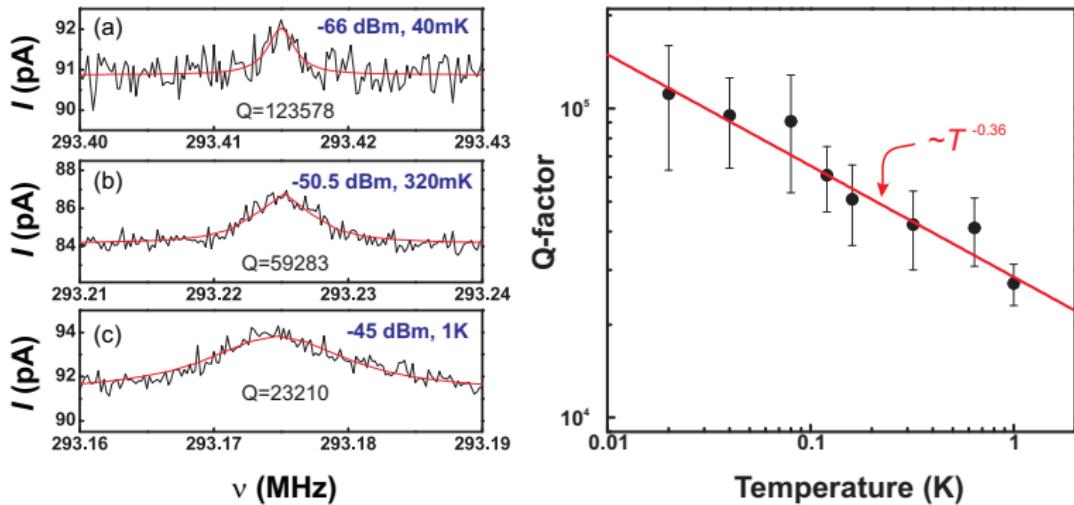


Driving into nonlinear response



- same temperature
- same working point V_g , V_{SD}
- low driving power:
symmetric, “linear” response
- high driving power:
asymmetric response, hysteresis
Duffing-like oscillator

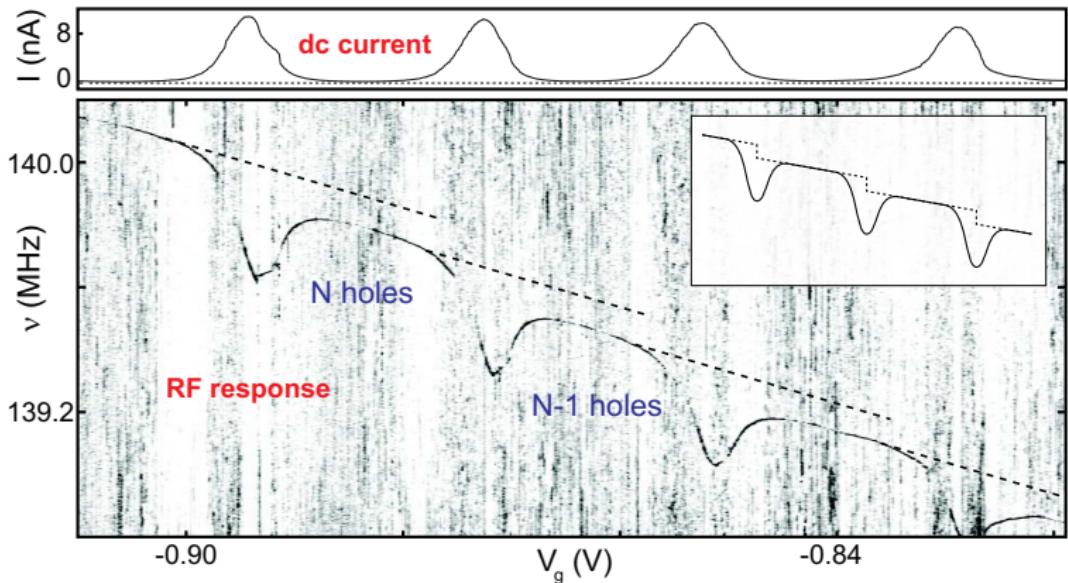
Temperature dependence of Q



$Q(T)$ fits power law prediction for intrinsic dissipation in nanotube

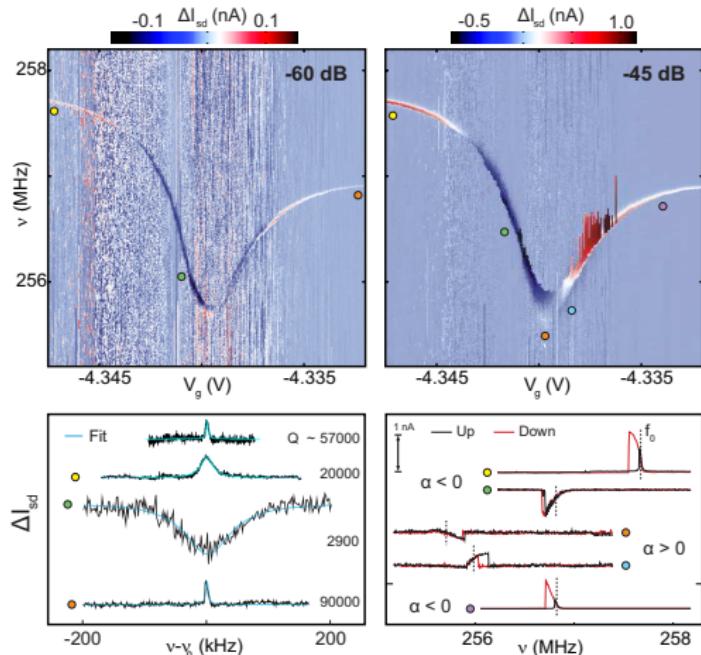
→ H. Jiang *et al.*, Phys. Rev. Lett. **93**, 185501 (2004)

Detailed $v(V_g)$: in SET, frequency decreases



“Coulomb blockade oscillations of **mechanical resonance frequency**”
electrostatic contribution to spring constant

Also Q and nonlinearity dominated by backaction

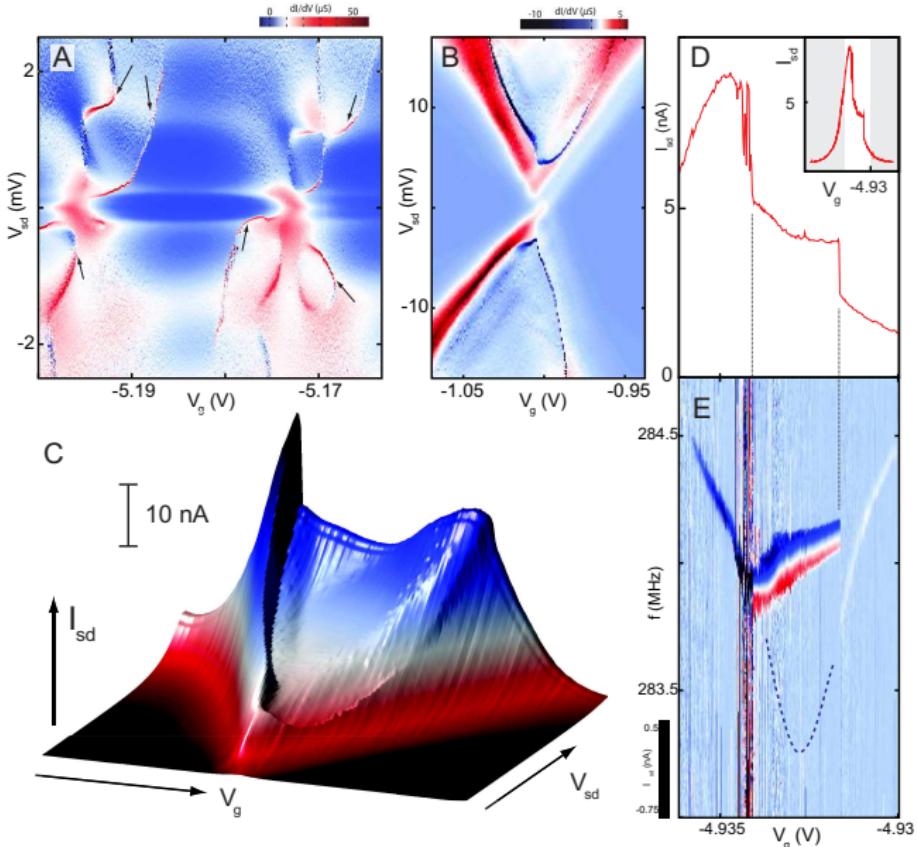


- Dissipation whenever charge can fluctuate
- Q decreases on SET peaks
- Nonlinearity dominated by tunneling
- Switches between weakening and softening spring

Summary, conclusion & outlook!

- $120 \text{ MHz} \lesssim v \lesssim 360 \text{ MHz}$, $Q \lesssim 150000$
- Self-detection of motion via dc current
- Easy driving into nonlinear oscillator regime
- $Q(T)$ is consistent with intrinsic dissipation model
- Single-electron steps of the resonance frequency
- Backaction of single electron tunneling on v , Q , nonlinearity
- **Self-excitation of motion!**
- Estimated motion amplitude at resonant driving $\sim 250 \text{ pm}$
compare thermal motion 6.5 pm , zero-point motion 1.9 pm
- Application as mass sensor: **sensitivity $4.2 \frac{\mu\text{m}}{\sqrt{\text{Hz}}}$**
- Without driving: **mechanical thermal occupation $n \simeq 1.2$**
- Stay tuned for more interesting results!!

Self-excitation of the resonator



Model for $v(V_g)$

- Electrostatic force between tube and backgate:

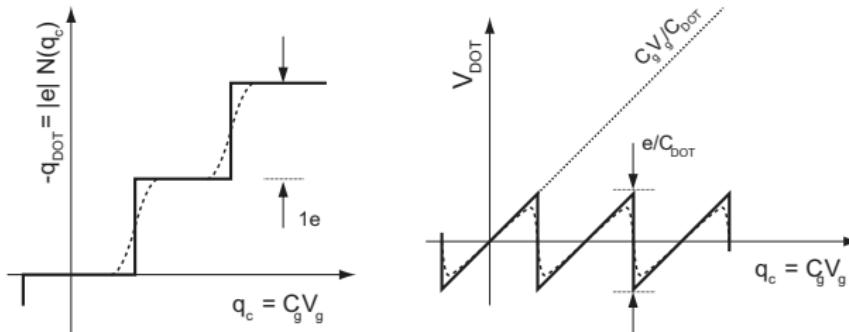
$$F_{\text{dot}} = \frac{1}{2} \frac{dC_g}{dz} (V_g - V_{\text{dot}})^2 \quad (1)$$

- Quantum dot voltage:

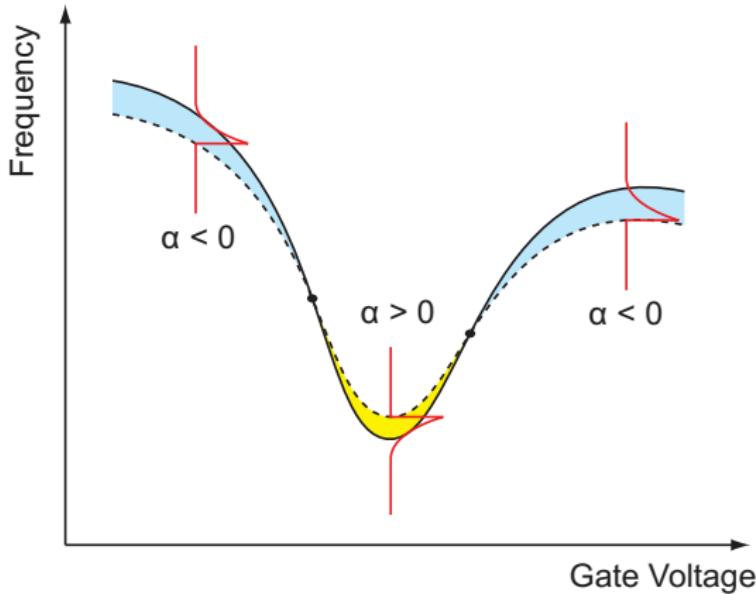
$$V_{\text{dot}} = \frac{C_g V_g + q_{\text{dot}}}{C_{\text{dot}}}, \quad q_{\text{dot}}(q_c) = -|e| \langle N \rangle(q_c), \quad q_c = C_g V_g \quad (2)$$

- Electrostatic contribution to spring constant:

$$k_{\text{dot}} = \frac{V_g(V_g - V_{\text{dot}})}{C_{\text{dot}}} \left(\frac{dC_g}{dz} \right)^2 \left(1 - |e| \frac{d\langle N \rangle}{dq_c} \right) \quad (3)$$



Model for $\alpha(V_g)$



$$\alpha_{\text{dot}} = -\frac{d^3 F}{dz^3} = \frac{d^2}{dz^2} k_{\text{dot}}(q_c) = V_g^2 \left(\frac{dC_g}{dz} \right)^2 \frac{d^2 k_{\text{dot}}}{dq_c^2} \quad (4)$$

The sign of α_{dot} follows the sign of the curvature of k_{dot} .